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## Double $\beta$ decay of $^{116}\text{Cd}$ . Final results of the Solotvina experiment and CAMEO project.

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Final results of the  $2\beta$  decay experiment, performed in the Solotvina Underground Laboratory with the help of low-background  $^{116}\text{CdWO}_4$  crystal scintillators, are presented. The revised half-life value of the two-neutrino  $2\beta$  decay of  $^{116}\text{Cd}$  has been measured as  $T_{1/2}^{2\nu} = 2.9_{-0.3}^{+0.4} \times 10^{19}$  yr, and the new half-life limit on the neutrinoless  $2\beta$  decay of  $^{116}\text{Cd}$  has been established as  $T_{1/2}^{0\nu} \geq 1.7 (2.6) \times 10^{23}$  yr at 90% (68%) C.L. The latter corresponds to a restriction on the Majorana neutrino mass:  $m_\nu \leq 1.7 (1.4)$  eV at 90% (68%) C.L. The CAMEO project could further advance sensitivity of the  $^{116}\text{Cd}$  experiment to  $m_\nu \leq 0.05$  eV.

Observations of neutrino oscillations manifest the non-zero neutrino mass and provide important motivation for highly sensitive experiments on neutrinoless double beta ( $0\nu 2\beta$ ) decay. However, this process still remains unobserved, and only half-life limits for  $0\nu 2\beta$  mode were obtained (see, e.g., reviews [1–3]). One of the most sensitive  $2\beta$  experiments has been performed in the Solotvina Underground Laboratory [4] by the Kiev-Firenze collaboration with the help of the enriched cadmium tungstate ( $^{116}\text{CdWO}_4$ ) crystal scintillators [5,6]. Here, we briefly describe the main results of this experiment and discuss the further advancement of its sensitivity.

The description of the set-up and its performance have been already published [5,6]. Four cadmium tungstate  $^{116}\text{CdWO}_4$  scintillators (enriched in  $^{116}\text{Cd}$  to 83%) with a total mass of 330 g were used. They were surrounded by an active shield made of 15  $\text{CdWO}_4$  crystals of large volume with a total mass of 20.6 kg. The whole  $\text{CdWO}_4$  array was situated within an additional active shield made of plastic scintillator  $40 \times 40 \times 95$  cm, thus, a complete  $4\pi$  active shield of the main ( $^{116}\text{CdWO}_4$ ) detector was provided. The outer passive shield consists of high-purity copper (3 – 6 cm), lead (22.5 – 30 cm) and polyethy-

lene (16 cm). For each event in the  $^{116}\text{CdWO}_4$  detector, the amplitude of a signal, its arrival time and pulse-shape (2048 channels with 50 ns width) were recorded. The energy resolution of the detector was  $\text{FWHM} = 8.0\%$  at 2615 keV. Due to active and passive shields, and as a result of the time-amplitude [7,8] and pulse-shape analysis [9,10] of the data, the background rate of the  $^{116}\text{CdWO}_4$  detector in the energy region 2.5 – 3.2 MeV ( $Q_{\beta\beta}$  of  $^{116}\text{Cd}$  is 2.8 MeV) was reduced down to 0.037(10) counts/(yr·kg·keV) [5,6]. It is the lowest background rate which has ever been reached with crystal scintillators.

The background spectrum measured during 12649 h after set-up upgrading in 1999 (the energy resolution and pulse-shape discrimination ability of the detector were improved) is shown in Fig. 1. Fit of the data gives the half-life value of the  $2\nu 2\beta$  decay of  $^{116}\text{Cd}$ :

$$T_{1/2}^{2\nu} = 2.9 \pm 0.06(\text{stat})_{-0.3}^{+0.4}(\text{syst}) \times 10^{19} \text{ yr.}$$

The main origins of the systematical uncertainties are possibly traces of the  $\beta$  active nuclides  $^{234m}\text{Pa}$  and  $^{90}\text{Y}$  (daughter of  $^{90}\text{Sr}$ ) in the  $^{116}\text{CdWO}_4$  crystals [6].

The experimental  $2\nu 2\beta$  Kurie plot (depicted in the inset of Fig. 1) is well described by the straight line in the 900 – 2500 keV energy region,

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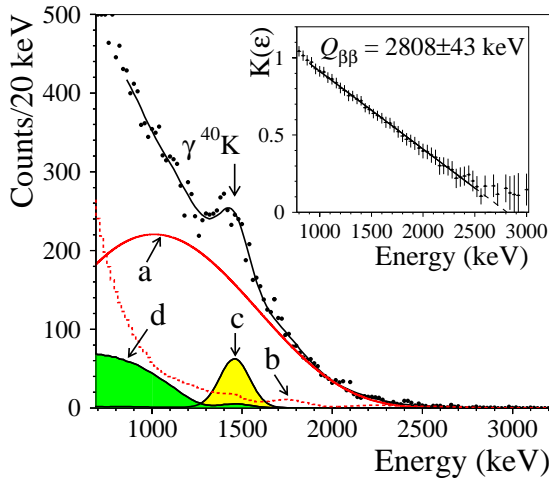


Figure 1. Spectrum of  $\gamma(\beta)$  events measured with  $^{116}\text{CdWO}_4$  detector (12649 h). Also shown are the most important model components: (a) the  $2\nu 2\beta$  spectrum of  $^{116}\text{Cd}$ ; (b) external  $\gamma$  background caused by radioactive contamination of the PMTs;  $^{40}\text{K}$  contamination of the (c) non-enriched and (d) enriched  $\text{CdWO}_4$  scintillators. Solid line represents the fit of the data in the 860 – 2700 keV energy interval. (Inset) The  $2\nu 2\beta$  decay Kurie plot and its fit.

and derived  $Q_{\beta\beta}$  energy ( $2808 \pm 43$  keV) is in accordance with the table value  $2805(4)$  keV. The obtained  $T_{1/2}^{2\nu}$  is in agreement with our preliminary result [5] and with those measured earlier in other experiments [7,11,12].

The part of the spectrum measured with the  $^{116}\text{CdWO}_4$  detector during 14183 h  $\{7.41 \times 10^{23}$  (nuclei of  $^{116}\text{Cd}) \times \text{yr}\}$  is presented in Fig. 2. Since any indications on the peak of  $0\nu 2\beta$  decay is absent in the experimental data, we set the bound on the half-life for  $0\nu 2\beta$  decay of  $^{116}\text{Cd}$ :

$$T_{1/2}^{0\nu} \geq 1.7 (2.6) \times 10^{23} \text{ yr at 90\% (68\%) C.L.}$$

Using this limit and calculations [13], we derive restrictions on the Majorana neutrino mass and right-handed admixtures in the weak interaction:  $m_\nu \leq 1.9$  eV,  $\eta \leq 2.5 \times 10^{-8}$ ,  $\lambda \leq 2.2 \times 10^{-6}$

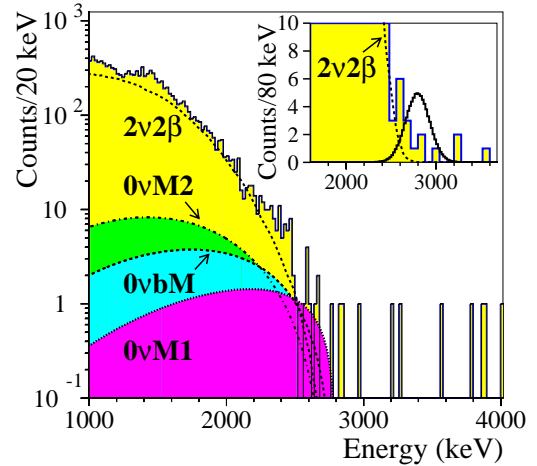


Figure 2. Part of experimental spectrum of the  $^{116}\text{CdWO}_4$ , together with the  $2\nu 2\beta$  distribution of  $^{116}\text{Cd}$ . Also are presented the excluded with 90% C.L. distributions of  $0\nu M1$ ,  $0\nu M2$  and  $0\nu$  decay with bulk Majoron emission with  $T_{1/2}^{M1} = 8 \times 10^{21}$  yr,  $T_{1/2}^{M2} = 8 \times 10^{20}$  yr, and  $T_{1/2}^M = 1.7 \times 10^{21}$  yr, respectively. In the inset the expected peak from  $0\nu 2\beta$  decay with  $T_{1/2} = 2 \times 10^{22}$  yr is shown.

at 90% C.L. Neglecting the right-handed contribution we get  $m_\nu \leq 1.7$  (1.4) eV at 90% (68%) C.L., and on the basis of [12,14] the limit is  $m_\nu \leq 1.5$  (1.2) eV. Besides, in accordance with ref. [15] the value of R-parity violating parameter of minimal SUSY standard model is restricted by our  $T_{1/2}^{0\nu}$  bound to  $\varepsilon \leq 7.0$  (6.3)  $\times 10^{-4}$  at 90% (68%) C.L.

To obtain limits for  $0\nu 2\beta$  decay with emission of one (M1), two (M2) and bulk [16] Majoron(s), the spectrum was analyzed in the energy region 1.6 – 2.8 MeV. Excluded with 90% C.L. distributions of  $0\nu 2\beta$  processes with Majoron emission are also shown in Fig. 2. In particular,  $T_{1/2}^{M1} \geq 0.8$  (1.8)  $\times 10^{22}$  yr at 90% (68%) C.L. Using this bound and calculations [17], the effective Majoron-neutrino coupling constant is restricted as  $g_M \leq 8.1$  (5.4)  $\times 10^{-5}$ , and on the basis of cal-

culatation [12,14] as  $g_M \leq 4.6 (3.1) \times 10^{-5}$ , which are among the strongest constraints reached in the direct  $2\beta$  decay experiments [3]. It should be stressed that all these restrictions were obtained with small  $^{116}\text{CdWO}_4$  detectors ( $\approx 0.3$  kg), in contrast with those for  $^{76}\text{Ge}$  studies ( $\approx 10$  kg of enriched HP  $^{76}\text{Ge}$  detectors [18]).

Therefore, the Solotvina experiment demonstrates that  $\text{CdWO}_4$  crystals possess several unique properties required for a  $2\beta$  decay experiment: low level of intrinsic radioactivity, good scintillation characteristics and pulse-shape discrimination ability, which allow one to reduce background effectively. But, to enhance the sensitivity to the neutrino mass bound to  $m_\nu \leq 0.05$  eV, we have to increase the mass of enriched  $^{116}\text{CdWO}_4$  detector and measuring time, to improve the energy resolution and to reduce the background of the detector further. Note that, for our  $^{116}\text{CdWO}_4$  detectors, the background in the 2.5–3.1 MeV energy interval is caused by the  $2\nu 2\beta$  decay of  $^{116}\text{Cd}$  ( $\approx 50\%$ ), internal  $^{232}\text{Th}$  contamination ( $\approx 30\%$ ), and by the external  $\gamma$  rays from the PMTs ( $\approx 20\%$ ). Hence, the further purification of the  $^{116}\text{CdWO}_4$  scintillators and the improvement of their energy resolution could additionally reduce the background. Namely, the CAMEO project [19], which intends to use  $\approx 100$  kg of  $^{116}\text{CdWO}_4$  crystals placed in a large volume of high purity liquid, calls also for improvement of the energy resolution (FWHM) at 2.8 MeV to 4% and for background reduction by factor  $\approx 100$ . The required radioactive contamination of  $^{116}\text{CdWO}_4$  crystals has to be less than  $\approx 10$   $\mu\text{Bq/kg}$ , both for  $^{228}\text{Th}$  and  $^{226}\text{Ra}$  [19]. Actually, even already existing  $\text{CdWO}_4$  crystals satisfy these requirements for  $^{226}\text{Ra}$  ( $< 4$   $\mu\text{Bq/kg}$ ) [6], but  $^{228}\text{Th}$  contamination varies in different crystals as (3 – 39)  $\mu\text{Bq/kg}$  [6,20]. At the same time, the required energy resolution FWHM = 4.3% at 2615 keV ( $^{232}\text{Th}$ ) has been already measured with the  $\text{CdWO}_4$  crystal  $\varnothing 4 \times 3$  cm placed in a liquid [19].

To conclude, fulfilment of the CAMEO experiment has practically no technical risk, and it seems to be very realistic because developed technique with the  $^{116}\text{CdWO}_4$  crystal scintillators is very simple and reliable. As the first step, the

CAMEO project is foreseen to produce  $\approx 10$  kg of  $^{116}\text{CdWO}_4$  crystals with required improved characteristics (radiopurity, energy resolution, etc.), which have been already obtained for several samples of  $\text{CdWO}_4$  scintillators. With such crystals, a half-life sensitivity  $\approx 10^{25}$  yr (corresponding to neutrino mass  $\approx 0.2$  eV) could be reached during 5 yr of measurements.

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